

The connection between the parsec-scale radio jet and γ -ray flares in the blazar 1156+295

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ABSTRACT

The blazar 1156+295 was active at γ -ray energies, exhibiting three prominent flares during the year 2010. Here, we present results using the combination of broad-band (X-ray through mm single-dish) monitoring data and radio-band imaging data at 43 GHz on the connection of γ -ray events to the ejections of superluminal components and other changes in the jet of 1156+295. The kinematics of the jet over the interval 2007.0–2012.5 using 43 GHz Very Long Baseline Array observations reveal the presence of four moving and one stationary component in the inner region of the blazar jet. The propagation of the third and fourth components in the jet corresponds closely in time to the active phase of the source in γ -rays. We briefly discuss the implications of the structural changes in the jet for the mechanism of γ -ray production during bright flares. To localize the γ -ray emission site in the blazar, we performed the correlation analysis between the 43 GHz radio core and the γ -ray light curve. The time lag obtained from the correlation constrains the γ -ray emitting region in the parsec-scale jet.

Key words: galaxies: active – galaxies: jets – quasars: individual: 1156+295 – gamma-rays: galaxies – radio continuum: galaxies.

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1 INTRODUCTION

Blazars are a subclass of active galactic nuclei (AGN) with a relativistic jet oriented close to the line of sight, which causes Doppler boosting of the jet emission and leads to strong variability at all wavebands from radio to γ -rays. It is generally accepted that the low-energy emission (from radio to UV or, in some cases, X-rays) is generated via synchrotron radiation by relativistic electrons in the jet plasma, while high-energy emission (from X-ray to γ -rays) is the result of inverse Compton (IC) scattering of seed photons by the same population of relativistic electrons. The seed photons could be either synchrotron photons generated in the jet (synchrotron self-Compton model; e.g. Atayan & Nahapetian 1989; Marscher 2014) or ambient photons (external Compton model; Begelman & Sikora 1987; Tavecchio et al. 2010). Several models have been proposed regarding the location of the γ -ray emission site relative to the central engine in blazars. Some of them constrain the location closer to the supermassive black hole (<0.1 – 1 pc), where the seed photons originate from the broad-line region (BLR) or the accretion disc (e.g. Tavecchio et al. 2010; Foschini et al. 2011; Rani et al. 2013). On the other hand, results from multifrequency studies suggest that the region where the bulk of the γ -rays is produced is usually located downstream of the canonical BLR (e.g. Marscher et al. 2010; Agudo et al. 2011; León-Tavares et al. 2011, 2012).

Many works have discussed the connection between the radio and γ -ray emission in blazars. The connection between γ -ray outbursts and radio flares, as well as structural changes observed in the jet with very long baseline interferometry (VLBI) were first established with data from the Energetic Gamma-ray Experiment Telescope (EGRET) detector on board the *Compton Gamma-ray Observatory* (Valtaoja & Teräsraanta 1996; Jorstad et al. 2001; Lähteenmäki & Valtaoja 2003) and, later, with the Large Area Telescope (LAT) on board the *Fermi Gamma-ray Space Telescope* (León-Tavares et al. 2011; Niippola et al. 2011). It was also found that blazars with strong γ -ray emission tend to be more luminous at radio frequencies (Kovalev et al. 2009) and exhibit highly superluminal motion, with the distribution of fastest speeds peaking at $\beta_{\text{app}} \sim 10c$ (Lister et al. 2009). Other studies have revealed details of the connection between low- and high-energy emission through extensive multifrequency variability studies of individual blazars (e.g. Marscher et al. 2010; Agudo et al. 2011; Schinzel et al. 2012; Wehrle et al. 2012; Jorstad et al. 2013; León-Tavares et al. 2013; Fuhrmann et al. 2014). Because blazars can exhibit a variety of behaviours when examined closely, it is important to carry out as many well-sampled multifrequency observational investigations as possible to sample the full range of behaviour and to identify common trends.

In this paper, we present results from a multifrequency study of the blazar 1156+295. This quasar, located at redshift $z = 0.729$ (Véron-Cetty & Véron 2010), displays strong variability across the electromagnetic spectrum. Prior to *Fermi*/LAT observations, the source was detected at γ -ray energies only in the second EGRET catalogue (Thompson et al. 1995). However, the *Fermi*/LAT with its better sensitivity had already detected 1156+295 at 1.6×10^{-7} photons $\text{cm}^{-2} \text{s}^{-1}$ after only three months of operation (Abdo et al. 2009). This object is classified as an optically violent variable and highly polarized quasar (cf. Wills et al. 1983, 1992; Fan et al. 2006). At radio frequencies, 1156+295 exhibits variability on both short and long time-scales (Hovatta et al. 2007; Savolainen & Kovalev 2008). On parsec to kiloparsec scales, the source exhibits a ‘core-jet’ structure. The components of the parsec-scale jet move at a wide range of apparently superluminal velocities, with component speeds up to $\sim 25c$ reported (Lister et al. 2013).

In 2010 August, a γ -ray flare with a flux ~ 10 times the average level was detected by the *Fermi*/LAT (Ciprini 2010). To study the flaring behaviour of the source and to determine the location of the γ -ray emission region, we perform a multiwavelength analysis. In Section 2, we describe the multifrequency data used in our analysis; in Section 3, we present results from Very Long Baseline Array (VLBA) observations and multifrequency light curves. In Section 4, we discuss the multifrequency connection and possible scenarios that can potentially explain the connection before drawing conclusions in Section 5.

We use a flat Λ cold dark matter cosmology with values, $H_0 = 68 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_m = 0.3$, and $\Omega_\Lambda = 0.7$ (Planck Collaboration 2013). This corresponds to a linear scale of 7.46 pc mas^{-1} at the redshift z of 0.729 for 1156+295 and a proper motion of 1 mas yr^{-1} corresponds to $42c$.

2 OBSERVATIONS AND DATA REDUCTION

2.1 γ -ray

The γ -ray fluxes over the energy range of 0.1–200 GeV were obtained by analysing the *Fermi*/LAT data from 2008 August 4 to 2011 December 31 using the *Fermi* Science Tools¹ version v9r33p0. To assure a high-quality selection of the data, an event class of 2 was applied with a further selection of zenith angle $>100^\circ$ to avoid contamination from photons coming from the Earth’s limb. The photons were extracted from a circular region centred on the source, within a radius of 15° . The instrument response functions P7REP_SOURCE_V15 were used (Ackermann et al. 2012).

We implemented an unbinned likelihood methodology using *glike* (Cash 1979; Mattox et al. 1996). This task models 31 point sources including our source within the region of interest (15°) obtained from the second *Fermi* Gamma-ray catalogue (hereafter 2FGL; Nolan et al. 2012). We fixed the model parameters of sources with significance $<2\sigma$ to the 2FGL value, while those of other sources were allowed to vary. We modelled our source using a simple power law. The Galactic diffuse emission and the isotropic background (sum of extragalactic diffuse and residual instrumental backgrounds) were also modelled at this stage, using the template ‘gll_iem_v05_rev1.fit’ and ‘iso_source_v05_rev1.txt’ – provided with the Science Tools. Our final fluxes were obtained from 7 d integrations, with a detection criterion such that the maximum-likelihood test statistic (TS; Mattox et al. 1996) exceeds 9 ($\sim 3\sigma$). For detections with $\text{TS} < 9$, 2σ upper limits were estimated using the profile likelihood method (Rolke, López & Conrad 2005).

2.2 X-ray and optical

In the X-rays, we obtained the *Swift* X-ray Telescope (XRT) data over the energy range, 0.3–10 keV, from an ongoing monitoring programme of *Fermi*/LAT monitored sources. The *Swift*/XRT data reduction method is discussed in Williamson et al. (2014). At optical wavelengths, we obtained *B*-, *V*-, *R*-, and *I*-band data from an ongoing monitoring programme of blazars at several observatories. The optical facilities include the Catalina Real-time Transient Survey (Drake et al. 2009),² Lowell Observatory (1.83 m Perkins Telescope equipped with the PRISM camera), Calar Alto (2.2 m

¹<http://fermi.gsfc.nasa.gov/ssc/data/analysis/documentation/Cicerone>

²<http://crts.caltech.edu/>

Telescope, observations under the MAPCAT³ programme), Liverpool 2 m Telescope, Crimean Astrophysical Observatory (0.7 m Telescope), and St. Petersburg State University (0.4 m Telescope). The optical data analysis procedures except for the Catalina data were performed as discussed in Jorstad et al. (2010).

2.3 Radio

The 230 GHz (1.3 mm) light curve was obtained at the Submillimeter Array (SMA). The source is included in an ongoing monitoring programme at the SMA to determine the fluxes of compact extragalactic radio sources that can be used as calibrators at mm wavelengths (Gurwell et al. 2007). Observations of the source are calibrated against known standards, typically Solar system objects (Titan, Uranus, Neptune, or Callisto). Data from this programme are updated regularly and are available at the SMA website.⁴

The 37 GHz single-dish fluxes were obtained from the observations made with the 13.7 m telescope at Aalto University Metsähovi Radio Observatory, Finland. The flux density scale is based on observations of the calibrator source DR 21, with NGC 7027, 3C 84, and 3C 274 used as secondary calibrators. A detailed description of the data reduction process and analysis is given in Teräsraanta et al. (1998).

To investigate the kinematics of the inner regions of the jet, we used 47 VLBA observations at 43 GHz from the Boston University blazar monitoring programme.⁵ The data reduction and calibration was performed as discussed in Jorstad et al. (2005).

We then modelled the complex visibility data with multiple components using the task *modelfit* in the *DIFMAP* program (Shepherd 1997), with each represented by a simple two-dimensional Gaussian brightness distribution. This method can identify components of the source structure that are closer than the resolution of the synthesized beam but are resolved by the longest baselines. Our model consisted of circular Gaussian components to parametrize the data in order to reduce the number of free parameters. The fit was considered to be good if the residual map rms noise was low and the reduced χ^2 statistic was ~ 1 . A fit to an additional component was deemed necessary only if it significantly improved the quality of the fit. No starting model was used during the model-fitting procedure. The uncertainties of the parameters of individual components were estimated with the *DIFWRAP* package (Lovell 2000), following the approach discussed in Rastorgueva et al. (2011). The model-fitting parameters are given in Table 1.

3 RESULTS

3.1 Multifrequency analysis

The light curves of the source 1156+295 from radio to γ -rays are presented in Fig. 1. At γ -ray energies, the average flux from 2008.6 to 2012 is 1.5×10^{-7} photons $\text{cm}^{-2} \text{s}^{-1}$. The source was in an active state from early 2010 to the beginning of 2011. Another period of prominent activity, but with lower intensity, occurred towards the end of 2011. Using the variability index discussed in Nolan et al. (2012),

$$\text{TS}_{\text{var}} = 2 \sum_i [\log L_i(F_i) - \log L_i(F_{\text{Const}})], \quad (1)$$

³<http://www.iaa.es/~iagudo/research/MAPCAT>

⁴<http://sma1.sma.hawaii.edu/callist/callist.html>

⁵<http://www.bu.edu/blazars/VLBAproject.html>

Table 1. 43 GHz model-fitting results. Columns are as follows: (1) observation epoch, (2) component identification (C0 and U correspond to core and unidentified component), (3) flux density in Jy, (4) distance from core in mas, (5) position angle with respect to core in degrees, (6) FWHM major axis of fitted Gaussian in mas. (This table is available in its entirety in the online journal.)

Epoch (1)	Component (2)	I (Jy) (3)	r (mas) (4)	P.A. (°) (5)	Maj. (mas) (6)
2007.45	C0	0.602	0.00	0.00	0.03
	U	0.068	0.22	8.80	0.11
2007.53	C0	0.509	0.00	0.00	0.03
	C1	0.078	0.14	350.7	0.16
2007.59	C0	0.350	0.00	0.00	0.05
	U	0.020	0.35	18.0	0.31
2007.83	C0	0.606	0.00	0.00	0.01
	C1	0.069	0.20	1.30	0.16
2008.04	C0	0.805	0.00	0.00	0.02
	C1	0.083	0.28	351.0	0.12
2008.16	C0	0.798	0.00	0.00	0.01
	U	0.244	0.06	19.2	0.19

where the value of the log likelihood in the null hypothesis, $\log L_i(F_{\text{Const}})$, corresponds to constant flux and those under the alternate hypothesis, $\log L_i(F_i)$, to variability, we found the source to be variable at the 99 per cent confidence level ($\text{TS}_{\text{var}} > 223.6$) with 177 degrees of freedom.

To characterize the active phase in γ -rays, we implemented a Bayesian blocks algorithm (Scargle et al. 2013) on the 2 d binned light curve (Fig. 2). This method generates a piecewise constant representation of the data by finding the optimal partition of the data. In turn, it globally optimizes the multiple change-point problem.⁶ Using a false positive rate⁷ of 0.01 and a prior value of the number of change points of 2.6, we obtained the block representation of the data presented in Fig. 2, which reveals four significant flares (A, B, C, and D). Each block in the figure is the weighted mean value of the observations within that block. We then obtained the parameters of each flare that are given in Table 2.

In the *Swift* 0.3–10 keV X-ray band, the source was in a high state around 2008.8, but due to the sparse sampling of the data, no further information could be inferred. At optical wavelengths, the source exhibited rapid variability on time-scales of days. During two major outbursts around 2008.4 and 2010.2, the brightness of the source increased by $\Delta R \sim 3$ mag within 2–3 months, while exhibiting variations on intraday time-scales during the rising phase of the flare. Such variations of 1156+295 have been previously reported by Raiteri et al. (1998) and Ghosh et al. (2000). The optical flare around 2010.2 occurred after component C3 (for more on jet components, see Section 3.2) was ejected. Unfortunately, during the γ -ray flares and after the ejection of C4, no optical observations were available owing to weather and seasonal visibility, hence prohibiting the multifrequency study in detail.

In the millimetre waveband (mm hereafter refers to 37 GHz), 1156+295 exhibits two flares with characteristic exponential rise

⁶In time series, a point at which a statistical model undergoes an abrupt transition, by one or more of its parameters jumping instantaneously to a new value, is called a ‘change point’.

⁷The probability of falsely reporting detection of a change point, similar to the value of alpha used in significance tests.

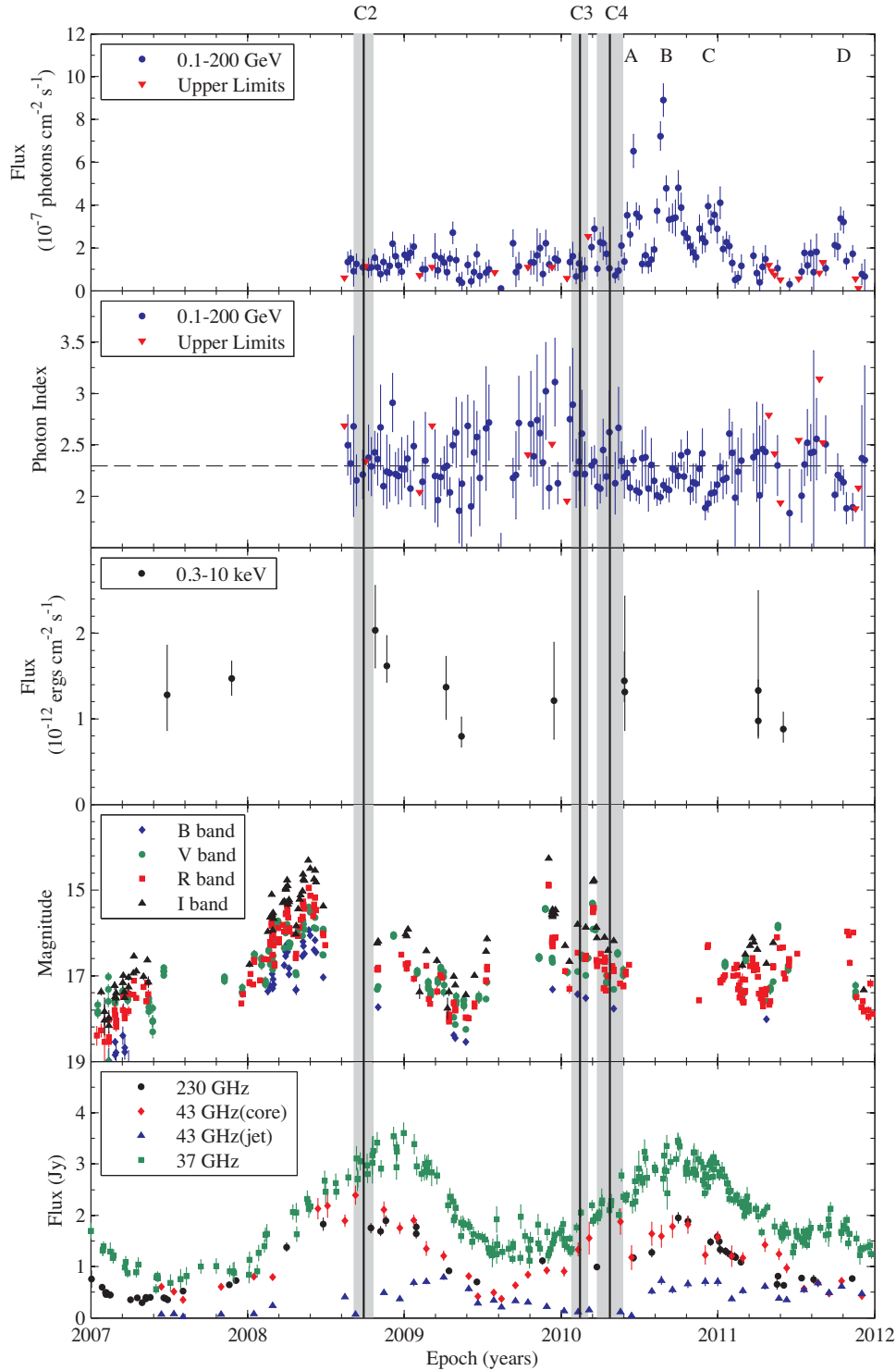


Figure 1. Light curves of 1156+295 from 2007 through 2011. From the top: (1) weekly binned γ -ray flux from *Fermi*/LAT at 0.1–200 GeV; 2σ upper limits are denoted as inverted red triangles. (2) Photon index of the weekly binned γ -ray light curve; the dashed line represents the 2FGL photon index estimate. (3) *Swift*/XRT X-ray flux at 0.3–10 keV. (4) Optical data at various bands. The gaps in the optical data are due to the proximity of the source to the Sun during certain annual intervals. (5) Variations at mm wavelengths in the bottom panel. The vertical lines are the ejection epochs of the components C2 (2008.74), C3 (2010.12), and C4 (2010.31) obtained from the VLBA data with their 1σ uncertainties denoted by the corresponding shaded interval.

and decay over the time period under study (Valtaoja et al. 1999). The first γ -ray flare occurs during the rising stage of the second mm flare, consistent with the analysis of mm and γ -ray light curves for a large sample of *Fermi*/LAT blazars presented in León-Tavares et al.

(2011). The 230 GHz SMA data are not considered in the following sections owing to their sparse sampling.

We associate the first mm flare with the ejection of component C2, while during the second mm flare two components (C3 and C4) were

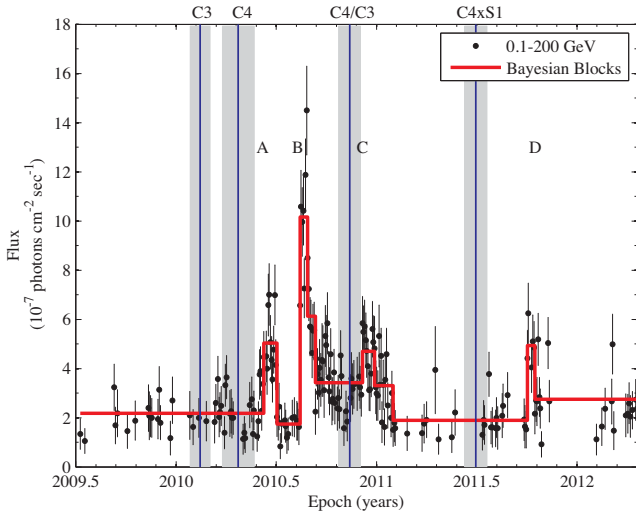


Figure 2. *Fermi* 2 d binned light curve at 0.1–200 GeV (filled black circles) using unbinned likelihood. Only measurements with $TS > 9$ are displayed here. Bayesian blocks representation is shown as a red line. The ejection epochs of the component C3 and C4 are plotted as vertical lines with 1σ uncertainties as shaded interval. C4/C3 and C4xS1 correspond to the time of splitting of component C4 from C3 and time of interaction of C4 with S1, respectively. See Section 4 for discussion.

ejected from the core. The rise time-scale for both flares is around 1.5 yr. The time to reach the quiescent state after flux maximum is around half a year for the first flare, while the second flare persists for a significantly longer period (around half a year) near the peak and then takes around a year to reach the quiescent state (Section 4). The broad peak and slower decay rate of the second mm flare, coupled with the slower apparent speed of C4 over the first segment of its trajectory, might be related to a lower magnetic field strength despite a higher density of electrons. This implies a possible deviation from the equipartition conditions during outbursts, as found previously by Homan et al. (2006). An alternative possibility is that both mm flares in reality consist of two (or more) individual flares, coming in rapid succession and blending together in the radio data. This would explain the unusual shape of both flares, which exhibit a more rapid decay than rise, unlike mm flares in general (Valtaoja et al. 1999; Hovatta et al. 2008).

3.2 Kinematics of the jet

In the kinematic analysis, the core is assumed to be stationary over the epochs. The cross-identification of the components in the subsequent epochs was based on the comparison of the parameters obtained from the Gaussian model fit. We have thus identified four

moving (C1, C2, C3, and C4) and one stationary (S1) component, based on the components' evolution in flux, distance, position angle, and size (see Table 1) derived from the 43 GHz VLBA data. Although C3 and C4 are blended for nearly one year after their ejections, the separation into two different components provides a smooth evolution of the jet features. The dynamics of the components, C3 and C4, were interpreted from the perspective of trailing shocks and forward/reverse structure for proper identification of the component (see Section 4). The total intensity images of the blazar at selected epochs with jet component locations marked according to modelling are shown in Fig. 3, while Fig. 4 shows the separation of the components from the core as a function of time.

The kinematics of the moving components were determined by fitting a polynomial using the method of least squares. We determined the order of the polynomial based on an F -test which tells whether addition of model parameters to fit the data is warranted by the level of misfit improvement. And thus, the first three moving components were fitted with a first-order polynomial while the best-fitting polynomial for component C4 was of second order based on an F -test with probability 1×10^{-5} . The fit yields the proper motion (μ) and the ejection epoch which is determined by back-extrapolating the fitted linear trajectory of every component (see Table 3). Following Jorstad et al. (2005), we obtained the accelerations for C4, both along and perpendicular to the jet to be 0.06 ± 0.01 and 0.34 ± 0.04 mas yr $^{-2}$. We estimated μ and t_o for C4 using first-order polynomial for the first seven epochs. No signs of acceleration were found in other components. We have calculated the apparent speed of every moving component (β_{app}) using the proper motion and luminosity distance (D_L) from the relation (Peebles 1993)

$$\beta_{\text{app}} = \mu \frac{D_L}{(1+z)}. \quad (2)$$

We estimate the physical parameters of the jet – Doppler factor, Lorentz factor, and viewing angle – under the assumption that the electrons emitting at 43 GHz lose energy mainly by radiative losses, so that the flux evolution is limited by the light-travel time across the component (see Table 3). The variability Doppler factor is thus estimated from the relation (Jorstad et al. 2005)

$$\delta_{\text{var}} = \frac{s D_L}{c \Delta t_{\text{var}} (1+z)}, \quad (3)$$

where s is the angular size of the component [i.e. the measured full width at half-maximum (FWHM) of the component multiplied by a factor of 1.8; cf. Pearson 1999] and Δt_{var} is the variability time-scale, defined as $\Delta t_{\text{var}} = dt / \ln(S_{\text{max}}/S_{\text{min}})$ (Burbidge, Jones & Odell 1974), where S_{max} and S_{min} are the measured maximum and minimum flux density of the component and dt is the time difference between S_{max} and S_{min} in years.

Table 2. Parameters of the γ -ray flares from Fig. 2. Columns are as follows: (1) flare ID, (2) duration of flare, (3) average flux over the duration, (4) time of peak flux, (5) peak flux of flare, (6) photon index at peak flux, (7) variability time-scale obtained using the same relation used for VLBA analysis and (8) time interval between flares.

Flare	Duration (d)	$\langle S_\gamma \rangle$ (10^{-7} photons cm^{-2} s^{-1})	T_γ^{peak} (4)	S_γ^{peak} (10^{-7} photons cm^{-2} s^{-1})	$\alpha_\gamma^{\text{peak}}$ (6)	Δt_{var} (d)	T_{diff} (d)
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)
A	24	4.78 ± 0.37	2010 June 19	7 ± 0.53	2.33 ± 0.45	11	–
B	16	9.95 ± 0.57	2010 August 26	14.48 ± 0.44	2.08 ± 0.13	15	68
C	24	4.71 ± 0.31	2010 December 6	5.84 ± 0.16	2.05 ± 0.19	19	102
D	12	4.46 ± 0.20	2011 October 4	6.24 ± 0.19	2.12 ± 0.33	8	302

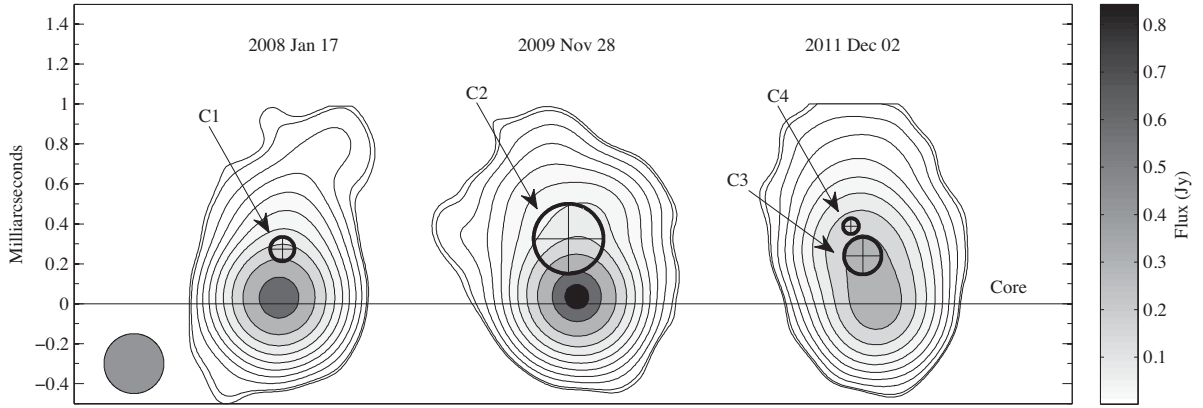


Figure 3. Structure of 1156+295 at selected epochs, with all four moving components identified. The images are convolved with a circular Gaussian beam of FWHM = 0.2 mas, as represented by the filled grey circle in the bottom-left corner.

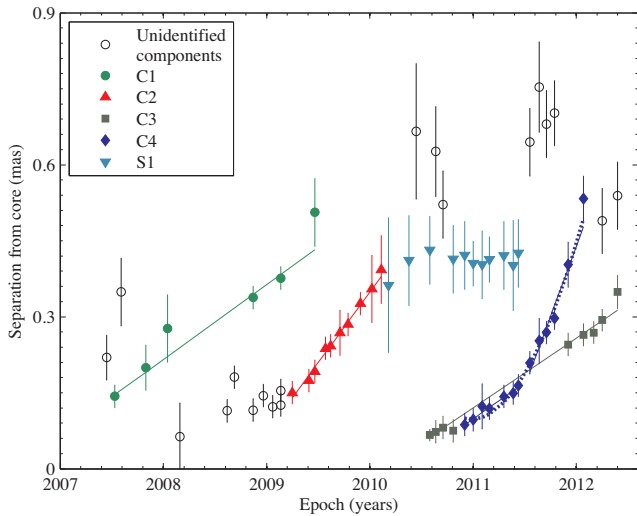


Figure 4. Separation of the components from the 43 GHz radio core versus time. Four moving components (C1, C2, C3, and C4) and one stationary component (S1) are identified in the jet. The solid lines indicate motion with no acceleration, while the dotted lines indicate accelerated motion fits. The displayed region corresponds to projected distances < 7 pc from the radio core. Different symbols indicate different components. Unidentified components are represented by unfilled circles.

By combining the variability Doppler factor with the apparent speed, we can calculate the variability Lorentz factor (Γ_{var}) and the viewing angle (θ_{var}) from the equations (Hovatta et al. 2009)

$$\Gamma_{\text{var}} = \frac{\beta_{\text{app}}^2 + \delta_{\text{var}}^2 + 1}{2\delta_{\text{var}}} \quad (4)$$

and

$$\theta_{\text{var}} = \arctan\left(\frac{2\beta_{\text{app}}}{\beta_{\text{app}}^2 + \delta_{\text{var}}^2 - 1}\right). \quad (5)$$

Using the average apparent speed, $\beta_{\text{app}} = 10.5c$, an upper limit to the jet viewing angle was obtained to be $\theta \leq 5.4^\circ$ from the relation: $\cos\theta = [\beta_{\text{app}}^2/(1 + \beta_{\text{app}}^2)]^{1/2}$. The average viewing angle from our results, $\theta_{\text{var}} = 3.5^\circ$, is consistent with the upper limit. The position angle of the moving components varies between -15° and $+30^\circ$ (consistent with the results reported in Jorstad et al. 2001). From the maximum viewing angle and the projected jet opening angle ($\phi_{\text{app}} \approx 45^\circ$), we constrained the maximum intrinsic jet opening angle, $\phi_{\text{int}} = \phi_{\text{app}} \sin\theta \lesssim 4.2^\circ$.

3.3 Localization of the γ -ray emission region

From the multifrequency light curves presented in Fig. 1, we see no activity in the γ -rays during 2008 when a major radio flare occurred in the same year. The flaring activity in the γ -rays during the

Table 3. Measured physical parameters of the components within 0.5 mas of the radio core. Columns are as follows: (1) component number, (2) number of epochs over which a component was identified, (3) proper motion, (4) apparent speed, (5) ejection epoch of the component, (6) variability time-scale, and (7–9) variability Doppler factor, Lorentz factor and viewing angle.

Component	Number of epochs	μ (mas yr ⁻¹)	β_{app} (c)	t_0 (yr)	Δt_{var} (yr)	δ_{var}	Γ_{var}	θ_{var} ($^\circ$)
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)
C1	6	0.147 ± 0.02	6.18 ± 0.8	2006.53 ± 0.12	0.75	10.85 ± 1.5	7.23 ± 1.4	4.56 ± 0.3
C2	10	0.278 ± 0.01	11.69 ± 1.5	2008.74 ± 0.06	0.44	18.54 ± 2.3	12.98 ± 2.1	2.79 ± 0.4
C3	9	0.137 ± 0.005	5.76 ± 0.4	2010.12 ± 0.05	1.06	7.9 ± 0.3	6.13 ± 1.2	6.83 ± 0.4
C4 ^a	7	0.142 ± 0.05	5.97 ± 0.8	2010.31 ± 0.08	0.55	15.37 ± 1.7	8.8 ± 0.5	2.52 ± 0.2
C4 ^b	6	0.552 ± 0.08	23.22 ± 2.3	$\sim 2011.4^c$	0.24	45.95 ± 2.1	28.85 ± 1.4	1 ± 0.08

Note: ^aEstimates obtained before acceleration.

^bEstimates obtained after acceleration.

^cTime of acceleration.

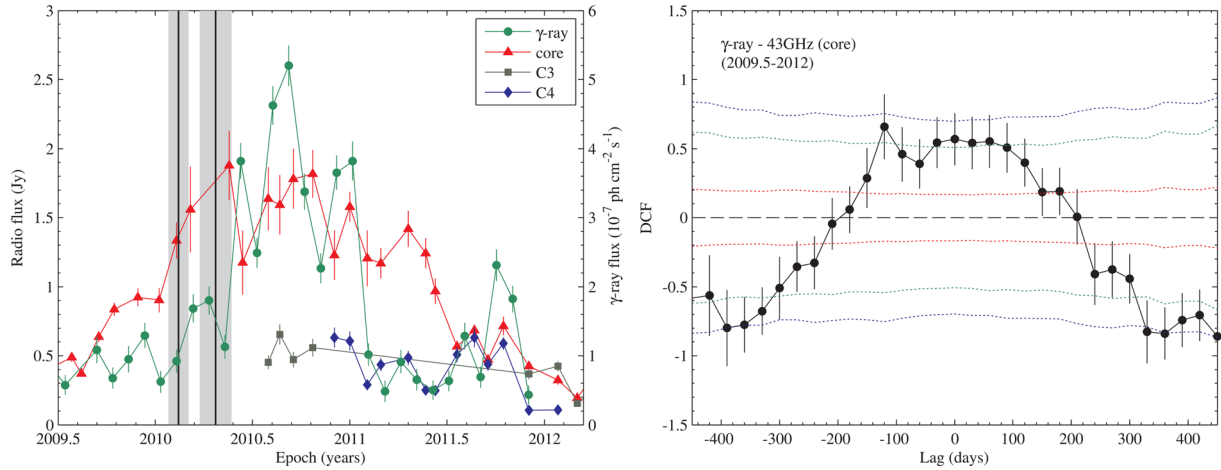


Figure 5. Left: VLBA core and components C3 and C4 light curves along with the monthly binned γ -ray light curve. Two vertical lines are the ejection epochs of the components C3 (2010.12) and C4 (2010.31). Right: cross-correlation between the monthly binned γ -ray and VLBA core light curve of 1156+295 for the time range -2009.5 – 2012 . Positive time lag indicates that activity in γ -rays precedes those in radio and the vice versa for the negative time lag. The significance levels of the correlations are denoted by red (1σ), green (2σ), and blue (3σ) dotted lines at positive (negative) DCF values estimated from Monte Carlo simulation as discussed in the text.

year 2010, however, corresponds closely in time to the variations in the VLBA core (see the left-hand panel in Fig. 5). Hence, to quantify this multifrequency behaviour, we perform the cross-correlation analysis using the discrete correlation function (DCF; Edelson & Krolik 1988) for unevenly sampled data. We applied the local normalization to the DCF, thus constraining it within the interval $[-1, +1]$ (White & Peterson 1994; Welsh 1999). We cross-correlated the VLBA core light curves with the monthly binned γ -ray light curves, to preserve the sampling. Only the time interval, 2009.5–2012, was considered.

The statistical significance of the cross-correlation is investigated using Monte Carlo simulations following Max-Moerbeck et al. (2013), under the assumption that the noise properties of the light curves can be described with a power-law power spectral density ($\propto 1/f^{-\alpha}$). The chosen power-law exponents are 1.5 for γ -rays (Abdo et al. 2010) and 2.3 for radio (Ramakrishnan et al., in preparation). We then simulated 5000 light curves using the power-law exponents with the method prescribed by Emmanoulopoulos, McHardy & Papadakis (2013). The simulated light curves are characterized to have the sampling pattern, mean, and variance as the observed light curves. In turn, we cross-correlated the simulated light curves using the DCF to estimate the distribution of random correlation coefficients at each time lag. We obtained 1σ , 2σ , and 3σ significance levels from the distribution. The result of the correlation analysis is summarized in Fig. 5.

The most prominent peak from our correlation analysis is located at a time lag of -120 d (~ 69 d in source frame) with >95.45 per cent significance, implying that the γ -rays are lagging the radio. We convert this time lag to linear distance travelled by the emission region, Δr , using the relation (Pushkarev, Kovalev & Lister 2010)

$$\Delta r = \frac{\beta_{\text{app}} c \Delta t_{\gamma, \text{radio}}^{\text{obs}}}{\sin \theta (1 + z)}, \quad (6)$$

where $\Delta t_{\gamma, \text{radio}}^{\text{obs}}$ is the observed time delay. Using the average apparent speed ($\beta_{\text{app}} = 7.4$) and average viewing angle ($\theta = 4.2^\circ$) obtained for all the components from Table 3, the location of the γ -ray emission region is constrained to lie at a projected distance ~ 6 pc from the 43 GHz core. This inference is in good agreement with the

analyses of other blazars by Jorstad et al. (2001), Lähtenmäki & Valtaoja (2003), Agudo et al. (2011), and León-Tavares et al. (2011). We also applied this method to the optical data, but owing to the presence of numerous gaps, no significant conclusion could be obtained.

We, however, note that the DCF peak being very broad, i.e. extending to also positive time lag (radio lagging), the result obtained above should be dealt with caution. The broadness of the DCF peak could be related to (i) the different time-scales of both events, i.e. the radio core light curve has a typical rise time of months while the γ -rays are significantly faster (rise time of days) (ii) and to the sparse sampling of the emission from the VLBI radio core.

3.4 Brightness temperature gradient along the jet

Abrupt changes in the brightness temperature (T_b) gradient can highlight regions in the jet where the density, magnetic field, or jet diameter changes rapidly. Hence, we have calculated T_b for each component from the equation (Kadler et al. 2004)

$$T_b = 1.22 \times 10^{12} \frac{S_{\text{comp}} (1 + z)}{d_{\text{comp}}^2 v^2}, \quad (7)$$

where S_{comp} is the component flux density in jansky and d_{comp} is the FWHM size of the circular Gaussian.

Brightness temperatures of all the model components in the jet in 1156+295 as a function of their distance from the core are plotted in Fig. 6. The variation in brightness temperature with distance is erratic for all components except C2. This behaviour is different from that expected for a stable conical jet with a straight axis and power-law dependences of the particle density, magnetic field strength, and the jet diameter on distance from the apex of the jet, r . The latter predicts that the brightness temperature along the jet can be described with a well-defined power-law index, f (Kadler et al. 2004), is

$$T_b \propto r^{-f}, \quad f = -l + n + b(1 + \alpha), \quad (8)$$

where α is the optically thin spectral index (flux density $S_\nu \propto \nu^{-\alpha}$), l , n , and b are power-law indices corresponding to the gradients of

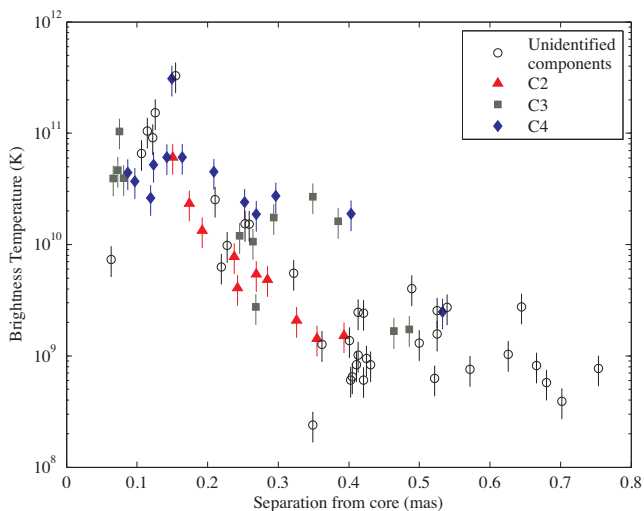


Figure 6. Brightness temperature of the components in the jet as a function of their distance from the core.

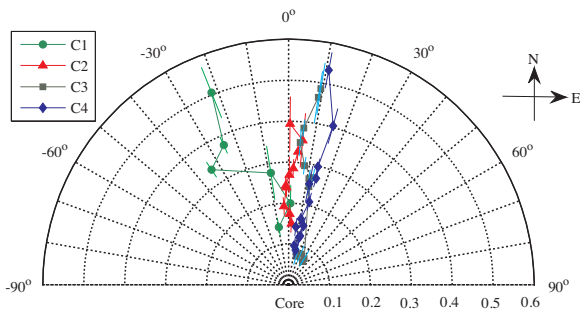


Figure 7. Propagation of the moving jet components with position angle. The separation of each moving component is shown by the semi-circle, the value of which is displayed at the bottom of the plot. The lines connecting the core to the circumference represent the position angle (marked in steps of 30°). The compass to the right shows the direction of the jet for clarification.

jet transverse size ($d \propto r^f$), power-law electron energy distribution ($n_e \propto r^{-n}$), and power-law magnetic field evolution ($B \propto r^{-b}$); if our line of sight subtends an angle $\lesssim (2\Gamma)^{-1}$ to the jet axis, $-l$ becomes $-2l$ in equation (8).

Most parsec-scale jets in AGN that do not show pronounced curvature do show a power-law decrease in brightness temperature with increasing distance from the core (e.g. Kadler et al. 2004; Pushkarev & Kovalev 2012; Schinzel et al. 2012). Component C2, with $f = 3.4$, shows such a behaviour, with mean position angle close to zero (Fig. 7). T_b for C3 and C4 is different and does not show a uniform decrease with the distance as in the case of C2. T_b of C4 increases at ~ 0.13 – 17 mas of the core and T_b of both C4 and C3 increases at ~ 0.3 – 0.4 mas from the core. This can be connected with splitting and interaction between C3 and C4 as well as possible interaction of C3 and C4 with S1.

4 DISCUSSION

The kinematics of the inner region in the jet of the blazar 1156+295, based on the 43 GHz VLBA observations, reveal the presence of four

moving, and one stationary, components. We find that the properties of the moving components differ from one another.

From a multifrequency perspective, the source was in an active state for almost a year in the γ rays and even longer at mm wavelengths (Figs 1 and 2). According to Fig. 2, the ejection of C3 and C4 corresponds to the beginning of strong γ -ray activity. The source was in an active state for more than three months (from flare B to C) in γ -rays before returning to a quiescent state. Flare B, with a peak flux of 1.4×10^{-6} photons $\text{cm}^{-2} \text{s}^{-1}$, is the brightest event observed in the γ -ray light curve. During the same time, we also notice an increase in flux of the component C3 by ~ 200 mJy (i.e. by 40 per cent; left-hand panel in Fig. 5).

Numerical hydrodynamical (HD) simulations of the dynamics of relativistic jets by Aloy et al. (2003) indicate that when the jet is perturbed at its injection point, the disturbance propagates downstream, spreading asymmetrically along the jet, and finally splitting into two regions. Both of these regions contain enhanced energy densities with respect to the underlying jet, and thus the synchrotron flux rises. The leading forward shock and trailing reverse shock have higher and lower Lorentz factors, respectively, than the underlying jet flow.

Another HD simulation finds that the interaction of the external medium with a strong shock pinches the surface of the jet, leading to the production of the trailing features (Agudo et al. 2001). These trailing shocks appear to be released in the wake of the primary superluminal component rather than ejected from the core. Hence, a single strong superluminal component ejection from the jet nozzle may lead to the production of multiple emission features through this mechanism.

From the context of the forward and reverse structures to the presence of a trailing shock, we discuss below possible models that might explain the activity in the jet during the evolution of the components C3 and C4 and its connection to the γ -ray activity.

(1) If the first four epochs of component C3 and the first seven epochs of component C4 in Fig. 4 correspond to the same physical disturbance in the jet, then, according to the discussion above, the combined feature could represent a forward/reverse structure, with the feature splitting around 2011.4 into two distinct components propagating at different speeds ($7.7c$ for C3 and $23.2c$ for C4 after the split). The Lorentz factors of the components are different (21.6 for C3 and 28.8 for C4 after separation), although not by as much as expected according to the simulations of Aloy et al. (2003). Forward shocks propagate faster than reverse shocks; when applied to 1156+295, this suggests that C4 could be a forward shock and C3 the corresponding reverse shock after 2011.4. However, the physical properties of C3 (flux, size, and position angle) during the first four epochs seem to be quite different from C4, hence casting doubt on the forward/reverse shock hypothesis.

(2) During the first seven epochs of C4, it could, by itself, represent a forward/reverse shock structure, since it displays significant variation of the flux (left-hand panel in Fig. 5) and size. It could then be regarded to split into two components around 2011.4, as in the first scenario. This hypothesis could explain the observation (Fig. 4) that, after splitting, the position angle of the two components remains the same (Fig. 7). However, it is difficult to reconcile this concept with the behaviour of the flux and size during the interaction of the moving component with stationary component S1 (see below).

(3) Component C4 can be classified as a trailing component, forming in the wake of the leading component, C3. Such a feature has been associated with the bright subluminal and superluminal jet

components in 3C 111 (Kadler et al. 2008) and 3C 120 (Gómez et al. 2001), as well as in 3C 273, 3C 345, CTA 102, and 3C 454.3 (Jorstad et al. 2005). After propagating over ~ 0.2 mas, C4 accelerates (Fig. 4), increasing the apparent speed to $23.2c$. This behaviour is in accordance with the simulation by Agudo et al. (2001), who find that the trailing components represent pinch waves excited by the main disturbance, so that an increase of their speed at larger distance reflects acceleration of the expanding jet. Under this scenario, the split of C4 from C3 towards the end of 2010 coincides with γ -ray flare C (denoted as C4/C3 in Fig. 2). Further investigation of this region in the jet is limited by the resolution of the VLBA.

The flux density evolution of the component C4 shows considerable variability which could be explained in terms of interaction with a stationary component (e.g. Gomez et al. 1997; León-Tavares et al. 2010), or by an increase in the Doppler boosting of the component at the positions where it is closer to the line of sight while travelling along a helical jet (e.g. Aloy et al. 2003; Hong et al. 2004). Although helicity has been studied in detail in this source (Hong et al. 2004; Zhao et al. 2011), from Fig. 4, the possibility of the interaction of C4 with S1 can be established. Component C4, after being accelerated around 2011.3, interacts with stationary component S1 (formed by early 2010) at 0.4 ± 0.04 mas around 2011.5 (denoted as C4xS1 in Fig. 2). Sub-flare D in the γ -rays occurs ~ 2 months after this interaction, which places the location of the sub-flare $\gtrsim 4$ pc (projected) from the radio core. During the interaction of C4 with S1, there is also an increase in the flux of C4 (~ 250 mJy), as shown in the left-hand panel of Fig. 5. Stationary components in this source have already been reported by Jorstad et al. (2001) and Zhao et al. (2011), but this is the first time that such a feature is found so close to the radio core in this source. Studies by Gomez et al. (1995, 1997) of relativistic HD and emission from jets show that when a moving component passes through a stationary feature produced by a standing shock, both components can appear to be blended into a single feature and the centroid of the merged components shifts downstream with respect to the pre-disturbance location of the stationary component. After the collision, the two components appear to split up, with the centroid of the quasi-stationary feature returning upstream. A similar model could be used to explain the feature exhibited by C4. The stationary feature could be produced at a bend in the jet, since this might also explain the observed position angle swing (Fig. 7).

The seemingly coincidence of the γ -ray activity with the radio core is clearly evident in the left-hand panel of Fig. 5. From the cross-correlation analysis, we were able to constrain the γ -ray emission site ~ 6 pc downstream of the radio core, which is consistent with previous results on other blazars (Marscher et al. 2010; Agudo et al. 2011; León-Tavares et al. 2011; Wehrle et al. 2012; Jorstad et al. 2013). However, before we can claim for the γ -ray emission region located further away from the radio core, we recall the caveat regarding the broad DCF peak mentioned in Section 3.3. We briefly enumerate various perspectives of the correlation analysis as follows.

(i) *Radio lagging the γ -rays*: except for the positive time lag between sub-flare C (around 2011) in the γ -rays and an increase in core flux during 2011.3, there is no clear evidence for the radio lagging the γ -rays (see the left-hand panel in Fig. 5).

(ii) *Almost zero time delay*: in these cases, when the same shock mentioned in (i) passes through the standing shock (aka the radio core), the radio emission starts to rise and peaks when the shock is at the centre of the radio core. Simultaneous γ -ray and radio outburst

could then be expected, if the source of seed photons is an external medium and the size of the radio core being very small or if the seed photons are from the jet due to shock–shock interaction. We can associate the flare and both sub-flares during 2010 in the γ -rays with a local maximum in the core fluxes. This connection supports the co-spatial origin of γ -ray and radio emission.

(iii) *Radio preceding the γ -rays*: it is evident in the light curve shown in Fig. 5 and as also discussed above that sub-flare D in the γ -rays corresponds to the local maximum of the component C4 in the jet which occurs from its interaction with the stationary component S1. This is similar to the results proposed by Agudo et al. (2011) for OJ 287 where the γ -ray emission was from the interaction of the moving shock with the quasi-stationary feature C1 located > 14 pc from the black hole. We refer to fig. 5 of Agudo et al. (2011) for a pictorial representation of the result that is also applicable here.

We do not have information on the true location of the radio core at 43 GHz. However, Pushkarev et al. (2012) find that the distance from the black hole to the radio core at 15 GHz is ~ 30 pc, well beyond the canonical BLR. Although the radio core at 43 GHz should be a factor of ~ 3 closer to the black hole, this still places it well beyond the inner parsec where the main BLR is expected to be located. The latter, in combination with our main finding that the γ -ray flare is produced after ~ 2 months of the start of component ejection, would allow us to rule out the model where the most intense γ -rays are produced by upscattering of photons from the BLR. However, recent results by León-Tavares et al. (2013) indicate that, in the quasar 3C 454.3, emission-line clouds can exist (and be ionized) at distances of several parsecs down the jet. This in turn suggests that IC scattering of line photons can occur even at distances well beyond the inner parsec. Our study constrains the γ -ray emission site to be close to the radio core during γ -ray flare B and farther downstream from the core during sub-flare D.

We could not obtain any significant correlations concerning the observed variability at other wavebands. The rapid optical variability could be produced by microflares in the accretion disc or through eclipsing of hotspots by the accretion disc (Wiita 1996). Some authors (e.g. Marscher et al. 2010; Agudo et al. 2011; Marscher 2014) have proposed that the rapid variations could be attributed to the presence of a turbulent magnetic field in the relativistic jets of blazars. However, owing to the lack of optical observations during the γ -ray flare, no further connection can be inferred.

No significant γ -ray event was found during the first mm flare, whereas the second mm flare was accompanied by strong γ -ray activity. This might be due to the presence of two components in the inner region of the jet during 2010 that contributed to variability of the jet emission through acceleration of the jet flow and interaction of components. Also, stationary components located downstream of the radio core have been found to play an important role in the release of energy (Arshakian et al. 2010; León-Tavares et al. 2010). No such feature was identified during the first mm flare. The radiative transfer modelling of the source by Aller et al. (2013) for the interval when the source was active in the γ -rays suggests that substantial part of the magnetic field energy density lies in an ordered component oriented along the jet axis from modelling the radio flare using four shocks. This implies that the γ -rays and the radio might be unrelated during the first mm flare (see appendix B in Nalewajko, Begelman & Sikora 2014). However, from the brightness temperature variations of components around 0.1–0.2 mas seen in Fig. 6, it is possible that a γ -ray flare could have occurred in 2008 prior to the start of the *Fermi* observations.

5 SUMMARY AND CONCLUSION

We have investigated the mm- γ -ray connection in the blazar 1156+295 by analysing a multiwavelength data set (briefly described in Section 2) and 43 GHz VLBA observations over the period of 2007–2012. Our findings are as follows.

(1) From the 43 GHz VLBI maps, we identify four moving and one stationary component (0.4 ± 0.04 mas from the radio core) with apparent speeds in the range 3–12 c shortly after ejection and viewing angles between 1° and 7° .

(2) In the γ -rays, one major flare and three sub-flares were noticed in 2010 and towards the end of 2011. The cross-correlation analysis to study the connection between the γ -ray activity and the radio core yielded a time lag of ~ 2 months in the source frame with the γ -rays lagging the radio. However, given the caveat mentioned in Section 3.3 regarding the correlation peak and its argument in Section 4, the flaring activity in γ -rays during 2010 can be associated with the radio core. The possibility of IC scattering of BLR photons might still be possible if an outflowing BLR surrounds the radio core.

(3) There is also evidence suggesting that the bulk of γ -rays was produced downstream of the radio core. This conclusion is suggested by the coincidence of sub-flare D in the γ -rays and the component interaction.

We have interpreted the complex changes in the parsec-scale structure of the jet from the perspective of forward/reverse shocks and trailing shocks. From consideration of the component evolution (Section 4), we judge the development of trailing shocks in the inner region of the jet to be the preferred scenario.

However, detailed modelling of shocks along with the polarization of the source at both optical and mm wavelengths could provide better constraints on all the physical parameters (Aller et al. 2014). Likewise, continued monitoring of the source and, if possible, higher frequency VLBI observations could improve our understanding of the jet structure and also help in localizing the emission regions of the high-energy flares.

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SUPPORTING INFORMATION

Additional Supporting Information may be found in the online version of this article:

Table 1. 43 GHz model-fitting results (<http://mnras.oxfordjournals.org/lookup/suppl/doi:10.1093/mnras/stu1873/-/DC1>).

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